

PROBABILITIES OF ROTATIONAL GAMMA TRANSITIONS OF TYPE E2 AND QUADRUPOLE MOMENTS OF DEFORMED NUCLEI WITH  $K = 1$  AND  $1/2$

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Expressions are given for  $B(E2)$  and  $Q$  for cases where these quantities depend not only on the direct matrix element, but also on the cross matrix element.

THE quadrupole moments  $Q$  of deformed nuclei and the reduced probabilities  $B$  for  $\gamma$  transitions of type E2 between levels of a rotational band may have anomalous values for the case where  $K = 1$  or  $1/2$  ( $K$  is the projection of the total angular momentum on the nuclear symmetry axis). For  $K = 1$  and  $1/2$  the quantities  $Q$  and  $B(E2)$  depend not only on the intrinsic quadrupole moment  $Q_0 \equiv \langle \chi_K | \hat{Q} | \chi_K \rangle$  but also on the cross matrix element  $\langle \chi_{-K} | \hat{Q} | \chi_K \rangle$  ( $\hat{Q}$  is the quadrupole moment operator,  $\chi_K$  is the function characterizing the internal state of the nucleus).

The anomaly in the values of  $Q$  and  $B(E2)$  in nuclei with  $K = 1$  and  $1/2$  is similar to the well-known anomaly in the magnetic moments  $\mu$  and reduced probabilities for rotational transitions,  $B(M1)$ , in nuclei with  $K = 1/2$ , and is due to the equivalence of positive and negative directions of the nuclear axis.

The quadrupole moment of a nucleus with angular momentum  $I$  and projection  $K = 1$  can be written as follows:

$$Q = Q_0 \frac{3 - I(I+1)[1 + (-1)^I 3b_0]}{(I+1)(2I+3)} \quad (1)$$

In particular, for the ground state of the rotational band where  $I = K = 1$ ,

$$Q = Q_0(1 + 6b_0)/10. \quad (2)$$

The coefficient  $b_0$  characterizes the ratio of the matrix elements  $\langle \chi_{-1} | \hat{Q} | \chi_1 \rangle / \langle \chi_1 | \hat{Q} | \chi_1 \rangle$ . The reduced probabilities for E2  $\gamma$  transitions between levels of a rotational band with  $K = 1$  can be expressed in terms of the same parameters  $Q_0$  and  $b_0$ . For transitions with  $I + 1 \rightarrow I$ ,

$$B(E2) = \frac{5}{16\pi} e^2 Q_0^2 (C_{I+1K}^{IK})^2 [1 - (-1)^{I-K} (I+1) b_0]^2, \quad (3)$$

For transitions with  $I + 2 \rightarrow I$ ,

$$B(E2) = \frac{5}{16\pi} e^2 Q_0^2 (C_{I+2K}^{IK})^2 [1 + (-1)^{I-K} b_0]^2, \quad (4)$$

where  $C \dots$  are Clebsch-Gordan coefficients.

In the case of  $K = 1/2$ , the quadrupole moment of the ground state of the rotational band  $I = K = 1/2$  is identically equal to zero. The expressions for the reduced probabilities  $B(E2)$  for transitions between levels of a rotational band with  $K = 1/2$  have the same form as (3) and (4), and differ only in the values of  $I$  and  $K$ . The coefficient  $b_0$  characterizes the ratio  $\langle \chi_{-1/2} | \hat{Q} | \chi_{1/2} \rangle / \langle \chi_{1/2} | \hat{Q} | \chi_{1/2} \rangle$ .

The magnitude and sign of the coefficient  $b_0$  are determined by the internal state of the nucleus. In particular cases it may turn out that the cross matrix element is small compared to  $Q_0$ . However, there is no basis for assuming that  $b_0 \ll 1$  for all nuclei. Therefore the measurement of a single value for  $B(E2)$  in nuclei with  $K = 1$  and  $1/2$  is insufficient for determining the intrinsic quadrupole moment and the deformation parameter of the nucleus.

There are a considerable number of deformed nuclei known at present which have states with  $K = 1$  or  $1/2$ . But the corresponding experimental data are available only for five nuclei with  $K = 1/2$ . These data are given in the table. There we also give the theoretical values calculated on the assumption that  $b_0 = 0$ . Agreement of experi-

Nucleus	Ratio	Theory for $b_0 = 0$	Experiment
Yb <sup>171</sup>	$\frac{B(E2, 1/2 \rightarrow 5/2)}{B(E2, 1/2 \rightarrow 3/2)}$	1.50	1.49 [1]
Tm <sup>169</sup>	$\frac{B(E2, 5/2 \rightarrow 3/2)}{B(E2, 5/2 \rightarrow 1/2)}$	0.28	0.31 [2,3]
W <sup>183</sup>	$\frac{B(E2, 5/2 \rightarrow 3/2)}{B(E2, 5/2 \rightarrow 1/2)}$	0.28	0.52 [4,5]
U <sup>235</sup>	$\frac{B(E2, 5/2 \rightarrow 3/2)}{B(E2, 5/2 \rightarrow 1/2)}$	0.28	0.16 [6,7]
Pu <sup>239</sup>	$\frac{B(E2, 5/2 \rightarrow 3/2)}{B(E2, 5/2 \rightarrow 1/2)}$	0.28	1.04 [8,9]

mental and theoretical values occurs only for  $\text{Yb}^{171}$  and  $\text{Tm}^{169}$ . For the other nuclei one observes deviations, i.e.,  $b_0$  and consequently also  $\langle \chi_{-K} | \hat{Q} | \chi_K \rangle$  are different from zero. However, the precision of the available experimental data is low, so that accurate measurements of  $Q$  and  $B(E2)$  for nuclei with  $K = 1$  and  $1/2$  would have considerable interest.

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<sup>3</sup>Kel'man, Metskhvarishvili, Preobrazhenskiĭ, Romanov and Tuchkevich, JETP 37, 639 (1959), Soviet Phys. JETP 10, 456 (1960).

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<sup>5</sup>Thulin, Rasmussen, Gallagher, Smith and Hollander, Phys. Rev. 104, 471 (1956).

<sup>6</sup>Freedman, Wagner, and Engelkemeir, Phys. Rev. 88, 1155 (1952).

<sup>7</sup>Gol'din, Novikova, and Kondrat'ev, Тезисы 8-го совещания по ядерной спектроскопии (Reports of the Eighth Conference on Nuclear Spectroscopy) Press, Academy of Sciences, U.S.S.R., 1958.

<sup>8</sup>J. O. Newton, Nucl. Phys. 5, 218 (1958).

<sup>9</sup>S. Baranov and K. Shlyagin, Атомная энергия 1, 52 (1956).

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