

indices make a small contribution relative to the parameter c^2 . Indeed, multiplication of n expressions $v + xE_i + yE_i'$ with identical indices i yields expressions of the same type, but with coefficients of the order of c^{2n} .

Taking into account the identity

$$\sum_{p=0}^{\infty} \frac{1}{n!} \sum_i \prod_{m=1}^p x_{i_m} = \exp \left\{ \sum_i x_i \right\},$$

the left side of which contains also terms with identical indices, we obtain from (19)

$$\begin{aligned} Z &= \sum_{(a)} A \exp \left\{ \sum_i (KE_i + v + xE_i + yE_i') \right\} \\ &= \sum_{(a)} \bar{A} \exp \left\{ \sum_i (K+x)E_i + \sum_i yE_i' \right\}. \end{aligned} \quad (20)$$

As seen from (20), in the principal approximation in c^2 the model under consideration is isomorphic to the Ising model with interaction along the diagonals. As we have seen above, two cases are possible in this Ising model: the isomorphism is either violated in the next-higher approximations, or is conserved in the next higher approximation and an assumption made Fisher^[7] and also by Anisimov, Voronel', and Gorodetskiĭ^[8] is satisfied, namely that the thermodynamic potential in the variable θ of a system with impurities is isomor-

phic to the corresponding thermodynamic potential of the pure substance.

Since the coefficients x and y in (20) are of the order of c^2 , the shift of the critical temperature due to the direct interaction is of the order of c^2 .

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ERRATA

Article by V. A. Belinskii, E. M. Lifshitz, and I. M. Khalatnikov, "The Oscillatory Mode of Approach to a Singularity in Homogeneous Cosmological Models with Rotating Axes" (**33**, 1061 (1971)).

In formula (A.7) for P_1^2 the last term in square brackets should be $+2\mu\nu\gamma_{13}\gamma_{23}$.

Article by V. S. Popov, "On the Properties of the Discrete Spectrum for Z Close to 137" (**33**, 665 (1971)).

1. The left side of formula 6 should read

$${}_x W_{k, ig}(x) / W_{k, ig}(x)$$

2. Formula (27') should read

$$\varepsilon_2(\alpha) = \begin{cases} 2^{-1/2} \left[1 + \frac{g}{2} \operatorname{ctg} gL \right] & \text{for } 0 < gL < \pi \\ g \operatorname{ctg} gL & \text{for } \pi < gL < 2\pi \end{cases}$$

Article by Yu. A. Bykovskii, N. N. Degtyarenko, V. F. Elesin, Yu. P. Kozyrev, and S. M. Sil'nov, "Mass Spectrometer Study of Laser Plasma" (**30**, 706 (1971)).

The system of equations (10) should read

$$\begin{aligned} I(z) &\approx \beta_1 \frac{W^{1/2}(\gamma' - 1)^{1/2}}{d^{2/3}} \ln [\beta_2 z^{1/2} (\gamma' - 1)^{1/2} W^{1/2}] \\ (\gamma' - 1)^{-1} &\approx \frac{3}{2} + \frac{Q(z)}{1 + z \beta_1 (\gamma' - 1)^{1/2} W^{1/2}} \end{aligned} \quad (10)$$

Article by Yu. N. Demkov and V. V. Ostrovskii, "n + l Filling Rule in the Periodic System and Focusing Potentials" (**35**, 66 (1972)).

On p. 67, Col. 1, line 2, in the phrase "the larger n at fixed N , the deeper the given level" n should be replaced by l . Correct formulation is implied in the remainder of the text. In the caption of Fig. 3 omit the last words "at the same instant of time." There are also slight errors in Fig. 1 for $Z = 41, 43-45, 55-56$, and $63-65$. In the right hand side of the formula for $f(\nu)$ (Appendix), the denominator should contain the factor $\Gamma(4l + \nu + 1)$ in place of $\Gamma(4l + n_r + 1)$.